

Waterjet

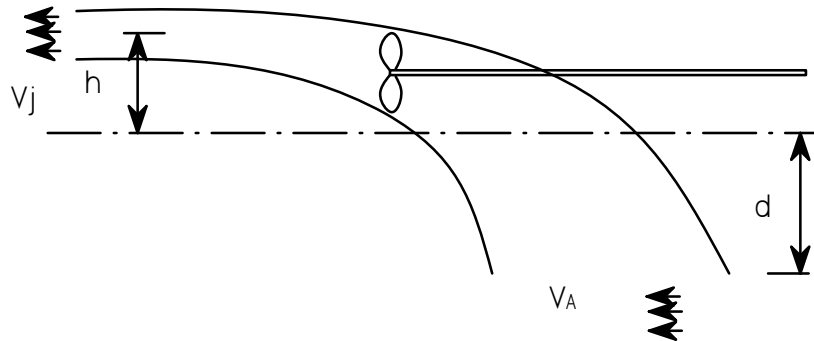
first draft 9/23/04 from Prof. Carmichael notes.
 9/17/06: modified to reflect w ($V \Rightarrow V_A$) and separate inlet
 and outlet pressure loss (in addition to drag) to reflect paper

- V_A velocity inlet
- w wake fraction
- V_s ship velocity
- V_j nozzle (outlet) velocity

$$V_A := V_s \cdot (1 - w)$$

$$T = \dot{m} \cdot (V_j - V_A)$$

$$\dot{m} = \text{mass_flow_rate}$$



at inlet centerline ... $P_{\text{local}} = P_{\text{atmos}} + \rho \cdot g \cdot d$

at this point ... total pressure (pitot tube) $P_{\text{oin}} = P_{\text{local}} + \frac{1}{2} \cdot \rho \cdot V_A^2 = P_{\text{atmos}} + \rho \cdot g \cdot d + \frac{1}{2} \cdot \rho \cdot V_A^2$

pressure at inlet to pump ... $P_{\text{op}} = P_{\text{oin}} - \rho \cdot g \cdot (d + h) = P_{\text{atmos}} - \rho \cdot g \cdot h + \frac{1}{2} \cdot \rho \cdot V_A^2$
 total pressure (pitot tube) ...

pressure at pump exit... total $P_{\text{oj}} = P_{\text{atmos}} + \frac{1}{2} \cdot \rho \cdot V_j^2$
 pressure (pitot tube) ..

total pressure increase $P_{\text{oj}} - P_{\text{op}} = P_{\text{atmos}} + \frac{1}{2} \cdot \rho \cdot V_j^2 - P_{\text{atmos}} + \rho \cdot g \cdot h - \frac{1}{2} \cdot \rho \cdot V_A^2 = \frac{1}{2} \cdot \rho \cdot (V_j^2 - V_A^2) + \rho \cdot g \cdot h$
 across pump ...

energy rise across the pump $\frac{P_{\text{oj}} - P_{\text{op}}}{\rho} = \frac{1}{2} \cdot (V_j^2 - V_A^2) + g h$
 per unit mass flow is ...

power absorbed by ideal pump is $P_{\text{pi}} = \dot{m} \cdot \frac{P_{\text{oj}} - P_{\text{op}}}{\rho} = \dot{m} \cdot \left[\frac{1}{2} \cdot (V_j^2 - V_A^2) + g h \right]$
 therefore ...

ideal efficiency is then ... and quasi propulsive coefficient is ...

$$\eta_i = \frac{P_{\text{Ti}}}{P_{\text{pi}}} \quad \eta_D = \frac{\text{effective_power}}{\text{power_delivered}} = \frac{P_E}{P_{\text{pi}}} = \frac{R \cdot V_s}{P_{\text{pi}}} = \frac{(1 - t) \cdot T \cdot V_s}{P_{\text{pi}}} = \frac{1 - t}{1 - w} \cdot \frac{T \cdot V_A}{P_{\text{pi}}} = \frac{1 - t}{1 - w} \cdot \frac{P_{\text{Ti}}}{P_{\text{pi}}} = \frac{1 - t}{1 - w} \cdot \eta_i$$

$$\eta_i = \frac{P_{\text{Ti}}}{P_{\text{pi}}} = \frac{T \cdot V_A}{P_{\text{pi}}} = \frac{\dot{m} \cdot V_A \cdot (V_j - V_A)}{\dot{m} \cdot \left[\frac{1}{2} \cdot (V_j^2 - V_A^2) + g h \right]} = \frac{2 \cdot V_A \cdot (V_j - V_A)}{V_j^2 - V_A^2 + 2g h} = \frac{2 \cdot V_A \cdot (V_j - V_A)}{V_j^2 - V_A^2 + 2g h} = \frac{2 \cdot \left(\frac{V_j}{V_A} - 1 \right)}{\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \frac{g h}{V^2}}$$

if $h = 0$

$$\eta_i = \frac{2 \cdot \left(\frac{V_j}{V_A} - 1 \right)}{\left(\frac{V_j}{V_A} \right)^2 - 1} = \frac{2}{\frac{V_j}{V_A} + 1}$$

same as propeller (we developed following in actuator disk

from actuator_disk.mcd using new variables to avoid duplication

$$\Delta v := V V_j - V V_A$$

$$V V := V V_A + \frac{\Delta v}{2}$$

$$\eta_I := \frac{T \cdot V V_A}{T \cdot V} \text{ simplify } \rightarrow 2 \cdot \frac{V V_A}{V V_A + V V_j} \quad \eta_I = \frac{2}{1 + \frac{V V_j}{V V_A}}$$

what are implications of $V V_j = V V_A$?

h cannot be negative (would be ducted prop, h limits efficiency)

Real waterjet with losses

net thrust of waterjet $T_{net} = T - \text{Drag}_{inlet} \quad T = m_{dot} \cdot (V_j - V_A)$

conventional drag coefficient $C_d = \frac{\text{Drag}}{\frac{1}{2} \cdot \rho \cdot v^2 \cdot A}$

drag coefficient of inlet $C_D = \frac{\text{Drag}}{\frac{1}{2} \cdot \rho \cdot v^2 \cdot A} = \frac{\text{Drag}}{\frac{1}{2} \cdot \rho \cdot V_A \cdot A \cdot V_A} = \frac{\text{Drag}}{\frac{1}{2} \cdot m_{dot} \cdot V_A}$

net thrust $T_{net} = T - \text{Drag}_{inlet} = m_{dot} \cdot (V_j - V_A) - C_D \cdot \frac{1}{2} \cdot m_{dot} \cdot V_A = m_{dot} \cdot V_A \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]$

net thrust power $P_{T_{net}} = T_{net} \cdot V_A = m_{dot} \cdot V_A^2 \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]$

delta p across pump must be increased to account for losses ...

we'll assume separate inlet and outlet losses

assume internal losses are ... $\sim 1/2 \cdot \rho \cdot v^2$

and the pump pressure rise is ...

$$\Delta p_{loss} = \Delta p_{in_loss} + \Delta p_{out_loss}$$

non-dimensional pressure loss coefficient is
and the real pump pressure rise is ...

$$K_{in} = \frac{\Delta p_{in_loss}}{\frac{1}{2} \cdot \rho \cdot V_A^2} \quad K_{out} = \frac{\Delta p_{out_loss}}{\frac{1}{2} \cdot \rho \cdot V_j^2}$$

$$P_{Oj} - P_{Op} = \frac{1}{2} \cdot \rho \cdot (V_j^2 - V_A^2) + \rho \cdot g \cdot h + \Delta p_{loss} = \frac{1}{2} \cdot \rho \cdot (V_j^2 - V_A^2) + \rho \cdot g \cdot h + K_{in} \cdot \frac{1}{2} \cdot \rho \cdot V_A^2 + K_{out} \cdot \frac{1}{2} \cdot \rho \cdot V_j^2$$

$$P_{Oj} - P_{Op} = \frac{1}{2} \cdot \rho \cdot V_A^2 \cdot \left[\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V_A^2} + K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2 \right]$$

ideal pump power is ...
$$P_{pi} = m_{dot} \cdot \frac{P_{oj} - P_{op}}{\rho} = m_{dot} \cdot \frac{1}{2} \cdot V_A^2 \cdot \left[\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V^2} + K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2 \right]$$

define η_p such that
$$\eta_p = \frac{P_{pi}}{P_p} \quad P_p = \text{actual_pump_power}$$

$$P_p = \frac{P_{pi}}{\eta_p} = \frac{m_{dot}}{\eta_p} \cdot \frac{P_{oj} - P_{op}}{\rho} = \frac{1}{2} \cdot \frac{m_{dot} \cdot V_A^2}{\eta_p} \cdot \left[\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V^2} + K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2 \right]$$

$$\eta_{real} = \frac{P_{T_net}}{P_p} = \frac{m_{dot} \cdot V_A^2 \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]}{\frac{1}{2} \cdot \frac{m_{dot} \cdot V_A^2}{\eta_p} \cdot \left[\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V_A^2} + \left[K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2 \right] \right]}$$

for a different form ... multiply numerator and denominator by $(V_A/V_j)^2$...

$$\eta_{real} = \frac{2 \cdot \eta_p \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]}{\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V_A^2} + K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2} = \frac{2 \cdot \eta_p \cdot \frac{V_A}{V_j} \cdot \left[\left(1 - \frac{V_A}{V_j} \right) - C_D \cdot \frac{1}{2} \cdot \frac{V_A}{V_j} \right]}{1 - \left(\frac{V_A}{V_j} \right)^2 + 2 \cdot \frac{g \cdot h}{V_j^2} + K_{in} \cdot \left(\frac{V_A}{V_j} \right)^2 + K_{out}}$$

and substitute μ for V_A/V_j ...
$$\eta_{real} = \frac{2 \cdot \eta_p \cdot \mu \cdot \left[(1 - \mu) - C_D \cdot \frac{1}{2} \cdot \mu \right]}{\mu^2 \cdot (K_{in} - 1) + 1 + K_{out} + 2 \cdot \frac{g \cdot h}{V_j^2}} = \frac{2 \cdot \eta_p \cdot \mu \cdot \left[(1 - \mu) - C_D \cdot \frac{1}{2} \cdot \mu \right]}{1 + K_{out} - \mu^2 \cdot (1 - K_{in}) + 2 \cdot \frac{g \cdot h}{V_j^2}}$$

and the quasi propulsive coefficient is then ..

$$\eta_D = \frac{1-t}{1-w} \cdot \eta_p \cdot \frac{2 \cdot \eta_p \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]}{\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V_A^2} + K_{in} + K_{out} \cdot \left(\frac{V_j}{V_A} \right)^2} = \frac{1-t}{1-w} \cdot \eta_p \cdot \frac{2 \cdot \mu \cdot \left[(1 - \mu) - C_D \cdot \frac{1}{2} \cdot \mu \right]}{1 + K_{out} - \mu^2 \cdot (1 - K_{in}) + 2 \cdot \frac{g \cdot h}{V_j^2}}$$

as from above ...

net thrust power
$$P_{T_net} = T_{net} \cdot V = m_{dot} \cdot V_A^2 \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]$$

first some comments to relate to previous lecture/notes version and Wärtzilä paper
 with $K_{out} = 0$ (N.B. this just means lumping all the pressure losses into a factor of $1/2 \cdot \rho \cdot V_A^2$ and
 accounting for a drag increase due to the inlet ...

$$\eta_D = \frac{1-t}{1-w} \cdot \eta_p \cdot \frac{2 \cdot \eta_p \cdot \left[\left(\frac{V_j}{V_A} - 1 \right) - C_D \cdot \frac{1}{2} \right]}{\left(\frac{V_j}{V_A} \right)^2 - 1 + 2 \cdot \frac{g \cdot h}{V_A^2} + K}$$

this is the form previously

and ... with $C_D = 0$ and assuming $h = 0$ (i.e. head loss is small compared to other terms ...

this is the form in the paper with

$$\eta_D = \frac{1-t}{1-w} \cdot \eta_p \cdot \frac{2 \cdot \mu \cdot (1-\mu)}{1 + K_{out} - \mu^2 \cdot (1 - K_{in})}$$


$$K_{out} = \phi = \text{nozzle_loss_coefficient}$$


$$K_{out} = \frac{\Delta p_{out_loss}}{\frac{1}{2} \cdot \rho \cdot V_j^2}$$

$$K_{in} = \varepsilon = \text{inlet_loss_coefficient}$$

$$K_{in} = \frac{\Delta p_{in_loss}}{\frac{1}{2} \cdot \rho \cdot V_A^2}$$

at this point, assuming K , C_D , and η_p are constant, could differentiate wrt V_j/V_A (or μ) and determine V_j/V_A for max propulsive coefficient, but minimum weight usually determines parameters.

 pump background (Wislicenus)

 example