

### 6.003: Signals and Systems

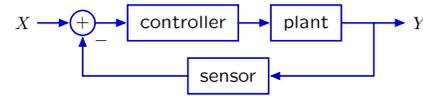
#### Discrete-Time Frequency Representations

April 13, 2010

#### Signals and/or Systems

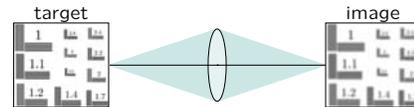
Two perspectives:

- feedback and control (focus on systems)



– Is the **system** stable?

- signal processing (focus on signals)



– Learn about target (**signal**) from the image (**signal**).

Fourier methods are especially useful in signal processing.

#### Historical Perspective

Broad range of CT signal-processing problems:

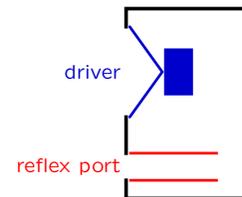
- audio
  - radio (noise/static reduction, automatic gain control, etc.)
  - telephone (equalizers, echo-suppression, etc.)
  - hi-fi (bass, treble, loudness, etc.)
- television (brightness, tint, etc.)
- radar and sonar (sensitivity, noise suppression, object detection)
- ...

Increasing important applications of DT signal processing:

- MP3
- JPEG
- MPEG
- MRI
- ...

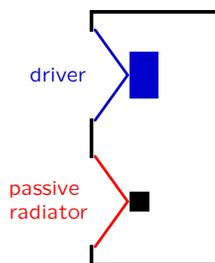
#### Signal Processing: Acoustical

Mechano-acoustic components to optimize frequency response of loudspeakers: e.g., “bass-reflex” system.



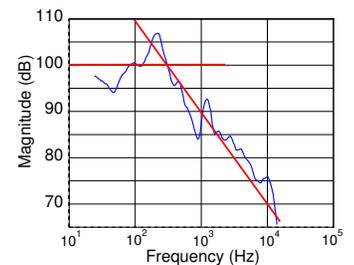
#### Signal Processing: Acoustico-Mechanical

Passive radiator for improved low-frequency performance.



#### Signal Processing: Electronic

The development of low-cost electronics enhanced our ability to alter the natural frequency responses of systems.

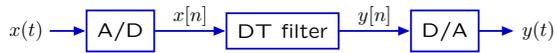


Eight drivers faced the wall; one pointed faced the listener.

Electronic “equalizer” compensates for limited frequency response.

**Signal Processing**

Modern audio systems process sounds digitally.



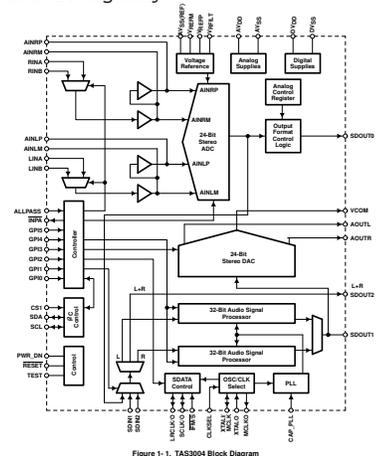
**Signal Processing**

Modern audio systems process sounds digitally.

Texas Instruments TAS3004

- 2 channels
- 24 bit ADC, 24 bit DAC
- 48 kHz sampling rate
- 100 MIPS
- \$7.70 (\$4.81 in bulk)

Courtesy of Texas Instruments.  
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**DT Fourier Series and Frequency Response**

Today: frequency representations for DT signals and systems.

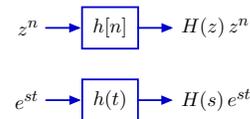
**Complex Geometric Sequences**

Complex geometric sequences are eigenfunctions of DT LTI systems.

Find response of DT LTI system ( $h[n]$ ) to input  $x[n] = z^n$ .

$$y[n] = (h * x)[n] = \sum_{k=-\infty}^{\infty} h[k]z^{n-k} = z^n \sum_{k=-\infty}^{\infty} h[k]z^{-k} = H(z)z^n.$$

Complex geometrics (DT): analogous to complex exponentials (CT)



**Rational System Functions**

A system described by a linear difference equation with constant coefficients → system function that is a ratio of polynomials in  $z$ .

Example:

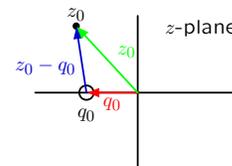
$$y[n - 2] + 3y[n - 1] + 4y[n] = 2x[n - 2] + 7x[n - 1] + 8x[n]$$

$$H(z) = \frac{2z^{-2} + 7z^{-1} + 8}{z^{-2} + 3z^{-1} + 4} = \frac{2 + 7z + 8z^2}{1 + 3z + 4z^2} \equiv \frac{N(z)}{D(z)}$$

**DT Vector Diagrams**

Factor the numerator and denominator of the system function to make poles and zeros explicit.

$$H(z) = K \frac{(z_0 - q_0)(z_0 - q_1)(z_0 - q_2) \cdots}{(z_0 - p_0)(z_0 - p_1)(z_0 - p_2) \cdots}$$



Each factor in the numerator/denominator corresponds to a vector from a zero/pole (here  $q_0$ ) to  $z_0$ , the point of interest in the  $z$ -plane.

Vector diagrams for DT are similar to those for CT.

**DT Vector Diagrams**

Value of  $H(z)$  at  $z = z_0$  can be determined by combining the contributions of the vectors associated with each of the poles and zeros.

$$H(z_0) = K \frac{(z_0 - q_0)(z_0 - q_1)(z_0 - q_2) \cdots}{(z_0 - p_0)(z_0 - p_1)(z_0 - p_2) \cdots}$$

The magnitude is determined by the product of the magnitudes.

$$|H(z_0)| = |K| \frac{|(z_0 - q_0)||z_0 - q_1||z_0 - q_2| \cdots}{|(z_0 - p_0)||z_0 - p_1||z_0 - p_2| \cdots}$$

The angle is determined by the sum of the angles.

$$\angle H(z_0) = \angle K + \angle(z_0 - q_0) + \angle(z_0 - q_1) + \cdots - \angle(z_0 - p_0) - \angle(z_0 - p_1) - \cdots$$

**DT Frequency Response**

Response to eternal sinusoids.

Let  $x[n] = \cos \Omega_0 n$  (for all time):

$$x[n] = \frac{1}{2} (e^{j\Omega_0 n} + e^{-j\Omega_0 n}) = \frac{1}{2} (z_0^n + z_1^n)$$

where  $z_0 = e^{j\Omega_0}$  and  $z_1 = e^{-j\Omega_0}$ .

The response to a sum is the sum of the responses:

$$\begin{aligned} y[n] &= \frac{1}{2} (H(z_0) z_0^n + H(z_1) z_1^n) \\ &= \frac{1}{2} (H(e^{j\Omega_0}) e^{j\Omega_0 n} + H(e^{-j\Omega_0}) e^{-j\Omega_0 n}) \end{aligned}$$

**Conjugate Symmetry**

For physical systems, the complex conjugate of  $H(e^{j\Omega})$  is  $H(e^{-j\Omega})$ .

The system function is the Z transform of the unit-sample response:

$$H(z) = \sum_{n=-\infty}^{\infty} h[n] z^{-n}$$

where  $h[n]$  is a real-valued function of  $n$  for physical systems.

$$H(e^{j\Omega}) = \sum_{n=-\infty}^{\infty} h[n] e^{-j\Omega n}$$

$$H(e^{-j\Omega}) = \sum_{n=-\infty}^{\infty} h[n] e^{j\Omega n} \equiv (H(e^{j\Omega}))^*$$

**DT Frequency Response**

Response to eternal sinusoids.

Let  $x[n] = \cos \Omega_0 n$  (for all time), which can be written as

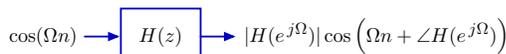
$$x[n] = \frac{1}{2} (e^{j\Omega_0 n} + e^{-j\Omega_0 n}) .$$

Then

$$\begin{aligned} y[n] &= \frac{1}{2} (H(e^{j\Omega_0}) e^{j\Omega_0 n} + H(e^{-j\Omega_0}) e^{-j\Omega_0 n}) \\ &= \text{Re} \{ H(e^{j\Omega_0}) e^{j\Omega_0 n} \} \\ &= \text{Re} \{ |H(e^{j\Omega_0})| e^{j\angle H(e^{j\Omega_0})} e^{j\Omega_0 n} \} \\ &= |H(e^{j\Omega_0})| \text{Re} \{ e^{j\Omega_0 n + j\angle H(e^{j\Omega_0})} \} \\ y[n] &= |H(e^{j\Omega_0})| \cos (\Omega_0 n + \angle H(e^{j\Omega_0})) \end{aligned}$$

**Frequency Response**

The magnitude and phase of the response of a system to an eternal cosine signal is the magnitude and phase of the system function evaluated on the unit circle.

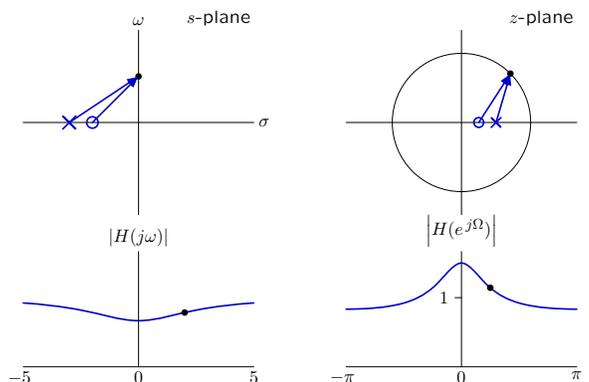


$$H(e^{j\Omega}) = H(z)|_{z=e^{j\Omega}}$$

**Comparison of CT and DT Frequency Responses**

CT frequency response:  $H(s)$  on the imaginary axis, i.e.,  $s = j\omega$ .

DT frequency response:  $H(z)$  on the unit circle, i.e.,  $z = e^{j\Omega}$ .



**Periodicity of DT Frequency Responses**

DT frequency responses are periodic functions of  $\Omega$ , with period  $2\pi$ .

If  $\Omega_2 = \Omega_1 + 2\pi k$  where  $k$  is an integer then

$$H(e^{j\Omega_2}) = H(e^{j(\Omega_1 + 2\pi k)}) = H(e^{j\Omega_1} e^{j2\pi k}) = H(e^{j\Omega_1})$$

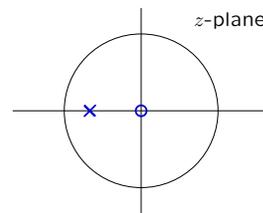
The periodicity of  $H(e^{j\Omega})$  results because  $H(e^{j\Omega})$  is a function of  $e^{j\Omega}$ , which is itself periodic in  $\Omega$ . Thus DT complex exponentials have many "aliases."

$$e^{j\Omega_2} = e^{j(\Omega_1 + 2\pi k)} = e^{j\Omega_1} e^{j2\pi k} = e^{j\Omega_1}$$

Because of this aliasing, there is a "highest" DT frequency:  $\Omega = \pi$ .

**Check Yourself**

What kind of filtering corresponds to the following?



- 1. high pass
- 2. low pass
- 3. band pass
- 4. band stop (notch)
- 5. none of above

**DT Fourier Series**

DT Fourier series represent DT signals in terms of the amplitudes and phases of harmonic components.

$$x[n] = \sum a_k e^{jk\Omega_0 n}$$

The period  $N$  of all harmonic components is the same.

**DT Fourier Series**

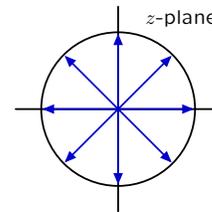
There are  $N$  distinct complex exponentials with period  $N$ .

If  $e^{j\Omega n}$  is periodic in  $N$  then

$$e^{j\Omega n} = e^{j\Omega(n+N)} = e^{j\Omega n} e^{j\Omega N}$$

and  $e^{j\Omega N}$  must be 1, and  $\Omega$  must be one of the  $N^{th}$  roots of 1.

Example:  $N = 8$

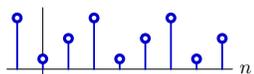


**DT Fourier Series**

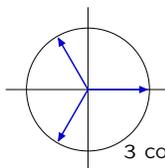
There are  $N$  distinct complex exponentials with period  $N$ .

These can be combined via Fourier series to produce periodic time signals with  $N$  independent samples.

Example: periodic in  $N=3$

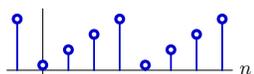


3 samples repeated in time

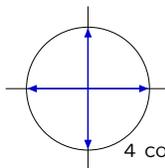


3 complex exponentials

Example: periodic in  $N=4$



4 samples repeated in time



4 complex exponentials

**DT Fourier Series**

DT Fourier series represent DT signals in terms of the amplitudes and phases of harmonic components.

$$x[n] = x[n + N] = \sum_{k=0}^{N-1} a_k e^{jk\Omega_0 n} \quad ; \quad \Omega_0 = \frac{2\pi}{N}$$

$N$  equations (one for each point in time  $n$ ) in  $N$  unknowns ( $a_k$ ).

Example:  $N = 4$

$$\begin{bmatrix} x[0] \\ x[1] \\ x[2] \\ x[3] \end{bmatrix} = \begin{bmatrix} e^{j\frac{2\pi}{4}0 \cdot 0} & e^{j\frac{2\pi}{4}1 \cdot 0} & e^{j\frac{2\pi}{4}2 \cdot 0} & e^{j\frac{2\pi}{4}3 \cdot 0} \\ e^{j\frac{2\pi}{4}0 \cdot 1} & e^{j\frac{2\pi}{4}1 \cdot 1} & e^{j\frac{2\pi}{4}2 \cdot 1} & e^{j\frac{2\pi}{4}3 \cdot 1} \\ e^{j\frac{2\pi}{4}0 \cdot 2} & e^{j\frac{2\pi}{4}1 \cdot 2} & e^{j\frac{2\pi}{4}2 \cdot 2} & e^{j\frac{2\pi}{4}3 \cdot 2} \\ e^{j\frac{2\pi}{4}0 \cdot 3} & e^{j\frac{2\pi}{4}1 \cdot 3} & e^{j\frac{2\pi}{4}2 \cdot 3} & e^{j\frac{2\pi}{4}3 \cdot 3} \end{bmatrix} \begin{bmatrix} a_0 \\ a_1 \\ a_2 \\ a_3 \end{bmatrix}$$

**DT Fourier Series**

DT Fourier series represent DT signals in terms of the amplitudes and phases of harmonic components.

$$x[n] = x[n + N] = \sum_{k=0}^{N-1} a_k e^{jk\Omega_0 n} \quad ; \quad \Omega_0 = \frac{2\pi}{N}$$

$N$  equations (one for each point in time  $n$ ) in  $n$  unknowns ( $a_k$ ).

Example:  $N = 4$

$$\begin{bmatrix} x[0] \\ x[1] \\ x[2] \\ x[3] \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & j & -1 & -j \\ 1 & -1 & 1 & -1 \\ 1 & -j & -1 & j \end{bmatrix} \begin{bmatrix} a_0 \\ a_1 \\ a_2 \\ a_3 \end{bmatrix}$$

**DT Fourier Series**

Solving these equations is simple because these complex exponentials are orthogonal to each other.

$$\begin{aligned} \sum_{n=0}^{N-1} e^{j\Omega_0 kn} e^{-j\Omega_0 ln} &= \sum_{n=0}^{N-1} e^{j\Omega_0 (k-l)n} \\ &= \begin{cases} N & ; k = l \\ \frac{1 - e^{j\Omega_0 (k-l)N}}{1 - e^{j\Omega_0 (k-l)}} = 0 & ; k \neq l \end{cases} \\ &= N\delta[k - l] \end{aligned}$$

**DT Fourier Series**

We can use the orthogonality property of these complex exponentials to sift out the Fourier series coefficients, one at a time.

$$\text{Assume } x[n] = \sum_{k=0}^{N-1} a_k e^{jk\Omega_0 n}$$

Multiply both sides by the complex conjugate of the  $l^{\text{th}}$  harmonic, and sum over time.

$$\begin{aligned} \sum_{n=0}^{N-1} x[n] e^{-jl\Omega_0 n} &= \sum_{n=0}^{N-1} \sum_{k=0}^{N-1} a_k e^{jk\Omega_0 n} e^{-jl\Omega_0 n} = \sum_{k=0}^{N-1} a_k \sum_{n=0}^{N-1} e^{jk\Omega_0 n} e^{-jl\Omega_0 n} \\ &= \sum_{k=0}^{N-1} a_k N\delta[k - l] = Na_l \end{aligned}$$

$$a_k = \frac{1}{N} \sum_{n=0}^{N-1} x[n] e^{-jk\Omega_0 n}$$

**DT Fourier Series**

Since both  $x[n]$  and  $a_k$  are periodic in  $N$ , the sums can be taken over any  $N$  successive indices.

Notation. If  $f[n]$  is periodic in  $N$ , then

$$\sum_{n=0}^{N-1} f[n] = \sum_{n=1}^N f[n] = \sum_{n=2}^{N+1} f[n] = \dots = \sum_{n=\langle N \rangle} f[n]$$

DT Fourier Series

$$a_k = a_{k+N} = \frac{1}{N} \sum_{n=\langle N \rangle} x[n] e^{-j\Omega_0 n} \quad ; \quad \Omega_0 = \frac{2\pi}{N} \quad (\text{"analysis" equation})$$

$$x[n] = x[n + N] = \sum_{k=\langle N \rangle} a_k e^{jk\Omega_0 n} \quad (\text{"synthesis" equation})$$

**DT Fourier Series**

DT Fourier series have simple matrix interpretations.

$$x[n] = x[n + 4] = \sum_{k=\langle 4 \rangle} a_k e^{jk\Omega_0 n} = \sum_{k=\langle 4 \rangle} a_k e^{jk\frac{2\pi}{4}n} = \sum_{k=\langle 4 \rangle} a_k j^{kn}$$

$$\begin{bmatrix} x[0] \\ x[1] \\ x[2] \\ x[3] \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & j & -1 & -j \\ 1 & -1 & 1 & -1 \\ 1 & -j & -1 & j \end{bmatrix} \begin{bmatrix} a_0 \\ a_1 \\ a_2 \\ a_3 \end{bmatrix}$$

$$a_k = a_{k+4} = \frac{1}{4} \sum_{n=\langle 4 \rangle} x[n] e^{-jk\Omega_0 n} = \frac{1}{4} \sum_{n=\langle 4 \rangle} e^{-jk\frac{2\pi}{4}n} = \frac{1}{4} \sum_{n=\langle 4 \rangle} x[n] j^{-kn}$$

$$\begin{bmatrix} a_0 \\ a_1 \\ a_2 \\ a_3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & -j & -1 & j \\ 1 & -1 & 1 & -1 \\ 1 & j & -1 & -j \end{bmatrix} \begin{bmatrix} x[0] \\ x[1] \\ x[2] \\ x[3] \end{bmatrix}$$

These matrices are inverses of each other.

**Discrete-Time Frequency Representations**

Similarities and differences between CT and DT.

DT frequency response

- vector diagrams (similar to CT)
- frequency response on unit circle in  $z$ -plane ( $j\omega$  axis in CT)

DT Fourier series

- represent signal as sum of harmonics (similar to CT)
- finite number of periodic harmonics (unlike CT)
- finite sum (unlike CT)

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